



Taylor, R.J.E., Li, G., Ivanov, P., Childs, D.T.D., Stevens, B.J., Harrison, B., Babazadeh, N., Ignatova, O. and Hogg, R.A. (2016) Mode Control in Photonic Crystal Surface Emitting Lasers Through In-Plane Feedback. In: 2016 International Semiconductor Laser Conference (ISLC), Kobe, Japan, 12-15 Sep 2016, ISBN 9784885523069.

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Mode Control in Photonic Crystal Surface Emitting Lasers Through In-Plane Feedback

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Abstract— In this paper, we show the effect of lateral external optical feedback on an all semiconductor photonic crystal surface emitting laser (PCSEL). Initially a PCSEL is grown and fabricated with a square lattice of triangles, the device is shown to operate electrically driven at room temperature under continuous wave condition. We investigate, theoretically and experimentally, the effect of lateral feedback on the performance of a photonic crystal lasers. Demonstrating a reduction in mode competition, and a modification to spatial mode distribution, opening routes to all electronic beam steering and divergence control.

Index Terms—Photonic Crystal Lasers, Semiconductor Lasers

I. INTRODUCTION

THERE have been significant recent developments in photonic crystal surface emitting lasers [1,2], with a particular focus on increasing output power through the development of individual lasers [3,4] through photonic crystal design and through coherent laser arrays [5]. Recently, triangular PC features (on a square lattice) have been exploited to realize watt-level output powers [6]. However, such PC structures significantly reduce the frequency difference between modes [7]. Closely spaced modes may lead to mode competition and multimode emission at high powers, both of which are deleterious for most practical laser applications. Utilizing distributed, varying phase feedback introduced through cleaved facets we demonstrate mode control in photonic crystal surface emitting lasers. We are able to show a reduction in threshold, a change in far field, and a switch from multimode to single mode emission.

PCSEL structures generally consist of a photonic crystal layer positioned above the active element within a p-i-n laser structure. Figure 1a shows a schematic representation of a typical PCSEL structure. Through careful design of the photonic crystal (PC) layer it is possible to control the lasing properties, in particular the beam shape [4], and beam polarization [8]. Beam steering has also been demonstrated [9].

PCSELs have previously mainly been realized through wafer fusion [2], however there has been a recent growing trend towards epitaxial overgrowth [10-15].

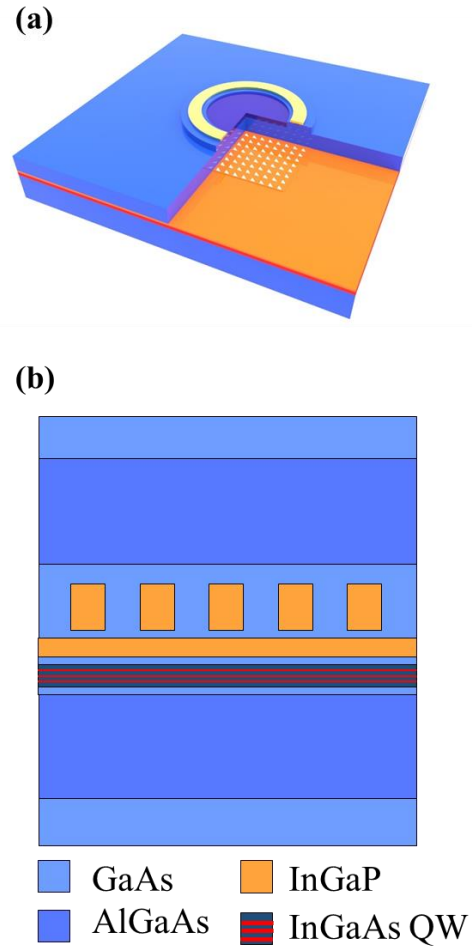


Figure 1 a) three dimensional rendered cut away schematic of PCSEL structure, b) two dimensional schematic of PCSEL layer structure.

“This work was financially supported by the Japan Society for the Promotion of Science (JSPS) under grant P15364 and the Engineering and Physical Sciences Research Council (EPSRC) under grant EP/K023195/1”.

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Gelleta et.al, [16] recently investigated simulated the effect of external feedback on PCSELS. With varying feedback intensity and phase to one side of a PCSEL (constant over the device side, but with varying values of phase shift), their model showed an influence on frequencies, threshold gain, gain margin, and field intensity. In this paper, we explore experimentally and theoretically the effect of external (in-plane) feedback on the lasing characteristics of all-semiconductor PCSELS. We begin by describing band-structure modelling which highlights the reduced mode frequency separation for triangular shaped PC features as compared to circular features. We then go on to describe a simple 1D model which predicts the effect of 1D feedback on the mode intensity within the PCSEL. We then go on to utilize a cleaved facet to introduce a distributed, varying facet phase across two orthogonal sides of the device ($\sim 4\text{-}5\pi$ distributed over ~ 450 periods ($150\text{ }\mu\text{m}$)). We describe the CW room temperature lasing characteristics of an all-semiconductor PCSEL and go on to discuss the effect of a single and second orthogonal cleave on the output characteristics of the device, and compare our results to theory.

II. DEVICE SIMULATION

In this section a photonic crystal is modelled by computing the definite frequency eigenstates of Maxwell's equations using fully vectorial plane wave expansion methods. The simulation is made using MIT photonic bands (MPB) [17], we consider the PC region to be infinite and to consist of InGaP and GaAs.

Figure 2 shows the TE photonic band diagram near the Γ point, and the inset shows an SEM micrograph of the PC (device realization is discussed in detail in the following section). The band structure is made up of two degenerate modes at lower frequencies and two non-degenerate (though closely spaced) modes to higher frequencies.

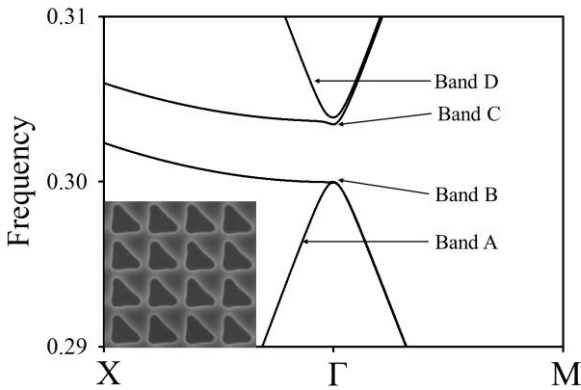


Figure 2 Simulated band structure of triangular PCSEL on a square lattice inset shows SEM of etched PCSEL region.

Figure 3 shows the dependence of Γ point band frequencies on the PC fill factor for (a) circular and (b) triangular atoms.

For both cases, as fill factor is increased, an increase in refractive index results in a monotonic reduction in mode frequencies.

As fill factor in circular PCs increases, we note a crossing of bands. This band crossing can result in a change from quadratic to linear dispersion at the Brillouin zone center and, if optimized, can lead to the realization of a laser with an accidental Dirac point. This is coincident with the degenerate/non-degenerate bands switching from the upper to the lower manifold. Accidental Dirac devices may exhibit many interesting phenomena [18-21].

For the PC with triangular features, we do not observe an intersection of bands. For the range where $f < 0.5$ (corresponding to triangular holes in-filled with another material of higher refractive index) the high-frequency bands of are non-degenerate but closely spaced. As the gain spectrum of the active element can be expected to span these two bands, competition between these closely spaced modes is observed. However, it is worth noting that Gelleta et.al. [17] demonstrated that lateral in-plane feedback can alter the threshold gain margin between photonic modes within the band structure.

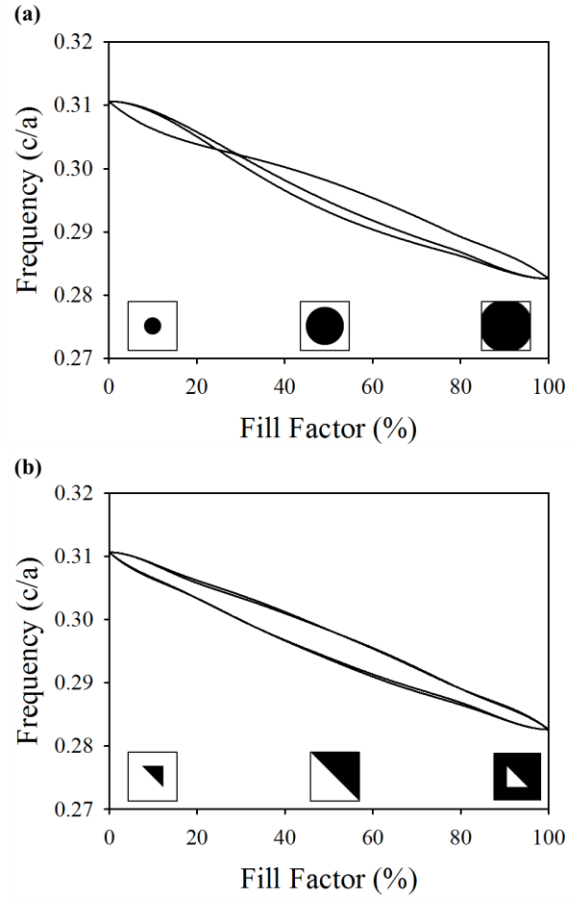


Figure 3 photonic crystal band frequencies as a function of fill factor for a) circular and b) triangular photonic crystals, insets show a schematic representation of photonic crystal fill factors of 10,50, and 90%.

In order to investigate the spatial distribution of optical power within the PCSEL we adopt a simple model which considers the PCSEL structure as a 1D waveguide. For a given column of unit cells within the 2D PC, orthogonal in-plane scattering into and out of the column is considered to be equal. Orthogonal out-of-plane scattering is considered equivalent to an

additional internal loss. A 1D model therefore provides a simple route to explore the effects of one dimensional feedback at the end of a column of PC unit cells.

The analysis is based on the transfer matrix method (TMM). Figure 4 shows a schematic representation of the layer structure simulated. The structure consists of a 1D periodically varying refractive index terminated by an air boundary, satisfying the piecewise homogeneous condition. Each period is characterized by a transfer matrix which represents the forward and backward travelling wave. An effective index method is used to calculate the vertical structural properties, where the effective index of each layer is calculated. A phase shift in external feedback is induced by varying the distance between the edge of the photonic crystal layer and the transparent semiconductor waveguide/air boundary. Threshold is obtained by analysis of the matrix elements when both real and imaginary part of M_{11} become zero.

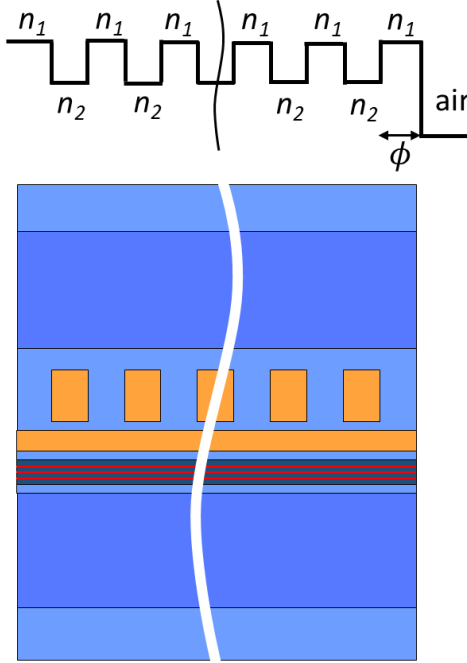


Figure 4 Schematic of simulated 1D PCSEL structure with refractive index

Figure 5 shows how the threshold gain changes as a function of reflection feedback phase for the fundamental (red) and second order mode (black). The phase shift is achieved in our model by varying the distance from photonic crystal region to the cleaved facet. The blue dashed line represents the threshold gain of an unclefted device (26 cm^{-1}).

For the fundamental mode threshold gain varies cyclically as the phase shift is varied from $-\pi$ to $+\pi$, the threshold gain reaches a minimum at $\frac{-\pi}{2}$ and a maximum at $\frac{\pi}{2}$. The threshold gain margin between the fundamental and second order mode is maximal for a facet phase of $+\pi$. As may be expected, facet phase plays a key role in determining whether in plane feedback will increase or decrease lasing threshold. It also suggests that threshold gain margin can be increased with suitable facet phase.

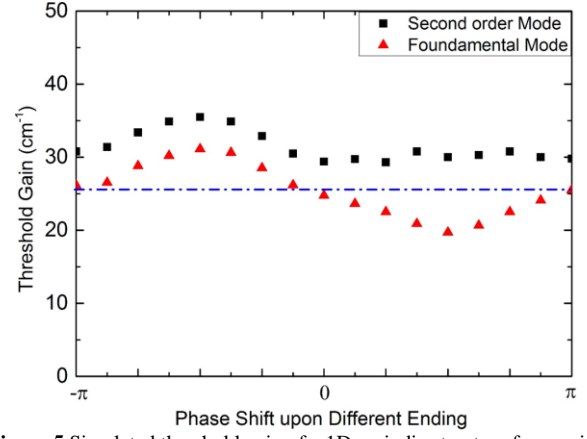


Figure 5 Simulated threshold gain of a 1D periodic structure for varying reflection phase for the fundamental (red) and second order mode (black).

Figure 6 shows the 1D simulation of E^2 field distribution for a $150 \mu\text{m}$ wide PCSEL. In figure 6a the boundary of the photonic crystal is transparent waveguide. The envelope of the E^2 field distribution has a single peak in the center of the photonic crystal region and a symmetric broad distribution. Figure 6b shows the E^2 field distribution when the waveguide is terminated with air, simulating a cleaved facet at the edge of the photonic crystal. The envelope of the E^2 field is observed to have a single peak which is shifted towards the cleave edge. The simulated structure in 6a with an air boundary layer shows significant narrowing of the envelope of the E^2 field distribution. This narrowing (by a factor of $2/3$) suggests the far-field will be broadened by a factor of $3/2$.

Figure 7 shows the calculated far-field pattern (1D) calculated from the near-field distribution in Fig 6, confirming a broadening of far-field angle of 1.5 times.

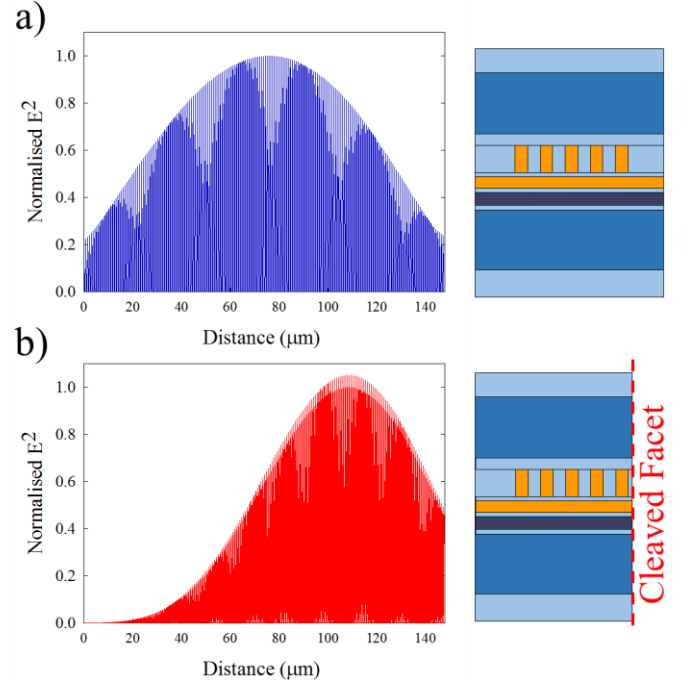


Figure 6 simulated E^2 field distribution of a 1D photonic crystal structure where the region around the photonic crystal is a) GaAs and b) air.

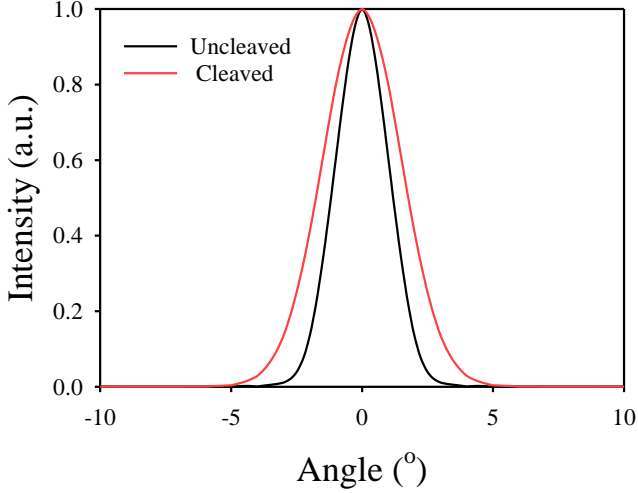


Figure 7 simulated far-field pattern for a 1D photonic crystal structure where the region around the photonic crystal is semiconductor (black) and air (red).

To summarize the simulation of feedback in 1D; The facet phase which produces the lowest threshold gain for the fundamental mode simultaneously produces the largest gain margin with the next lasing mode. Spatially, the introduction of feedback moves the center of the mode towards the origin of the feedback, and narrows the near-field envelope. These observations suggest that lateral feedback may not only be of importance in mode control, but also allow future possible routes to all electronic beam steering and divergence control.

III. DEVICE REALIZATION

Devices were grown by MOVPE on GaAs substrates cut 3° of (110). Initial growth consists of $1.5 \mu\text{m}$ n-type $\text{Al}_{0.4}\text{Ga}_{0.6}\text{As}$ lower cladding layer, 3 quantum well active region (8 nm $\text{In}_{0.2}\text{Ga}_{0.8}\text{As}$ QW's separated by 20 nm GaAs layers), a 40 nm p- $\text{In}_{0.48}\text{Ga}_{0.52}$ etch stop layer and a 150 nm $\text{In}_{0.48}\text{Ga}_{0.52}\text{P}$ layer (layer forms the PC region). The InGaP layer is patterned using electron beam lithography, where by right angled triangular holes are patterned into a SiO_2 hard mask. The pattern is transferred into the InGaP with a $\text{CH}_4/\text{H}_2/\text{O}_2$ reactive ion etch. The complete PC area consists of $150 \mu\text{m} \times 150 \mu\text{m}$ square area. Right angled triangles were chosen due to their apparent optimization of output power [6]. After etching, the wafer is cleaned using HF and placed back in the reactor. An overgrowth consisting of GaAs (to infill the holes and form the PC), a $1.5 \mu\text{m}$ p-type $\text{Al}_{0.4}\text{Ga}_{0.6}\text{As}$ upper cladding layer and finally a 400 nm p⁺ GaAs cap layer. The growth parameters for this structure have been optimized elsewhere [11]. After over growth, devices were fabricated by etching a 100 nm diameter mesa through the p⁺ cap layer above the PC region, finally contacts were defined to provide a $52 \mu\text{m}$ aperture. Figure 1a shows a 3D rendered cut away schematic representation of the PCSEL structure. Figure 1b shows a 1D schematic of the layer structure.

The completed structure consists of (from bottom to top) $1.5 \mu\text{m}$ n-type $\text{Al}_{0.4}\text{Ga}_{0.6}\text{As}$ lower cladding layer, 3 quantum well

active region (8 nm $\text{In}_{0.2}\text{Ga}_{0.8}\text{As}$ QW's separated by 20 nm GaAs layers), a 40 nm p- $\text{In}_{0.48}\text{Ga}_{0.52}$ etch stop layer, 150 nm photonic crystal layer, p-type $1.5 \mu\text{m}$ $\text{Al}_{0.4}\text{Ga}_{0.6}\text{As}$ cladding layer, finally a 400 nm p⁺-type GaAs cap.

Cleaved facets are placed orthogonally along the edge of the photonic crystal region, providing lateral in-plane feedback with spatially varying feedback phase. Figure 8 shows a schematic of the photonic crystal lattice geometry indicating misalignment between PCSEL plane and cleave plane, in this device the misalignment provides a feedback phase variation of $\sim 4\text{--}5\pi$ over the ~ 450 photonic crystal periods ($150 \mu\text{m}$). Utilizing a distributed, varying feedback phase is expected to select the lowest threshold 2D mode due to the 2D nature of the PC feedback.

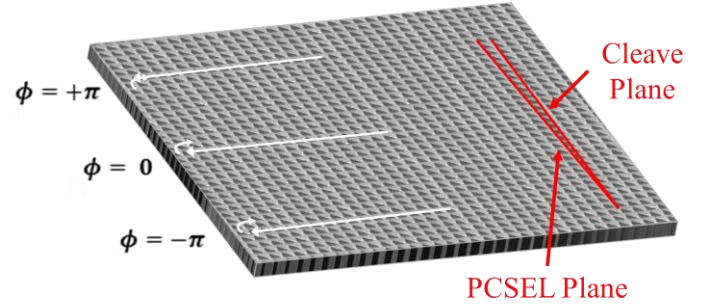


Figure 8 schematic showing photonic crystal lattice geometry indicating misalignment between PCSEL plane and cleave plane, showing in plane facet reflection phase.

IV. EXPERIMENTAL RESULTS

Figure 9b shows the CW room temperature LI of a PCSEL device, the device threshold is 112 mA ($J=1.43 \text{ kA}\cdot\text{cm}^{-2}$). The inset shows the far-field pattern which is symmetric with 1.2° divergence. Figure 9a plots the lasing spectra of the virgin device for current ranging from 100 to 200 mA. Two lasing modes at 963 nm and 963.5 nm are observed, with a total linewidth of ~ 1.5 nm. The two distinct modes are attributed to modes C and D in Figure 2.

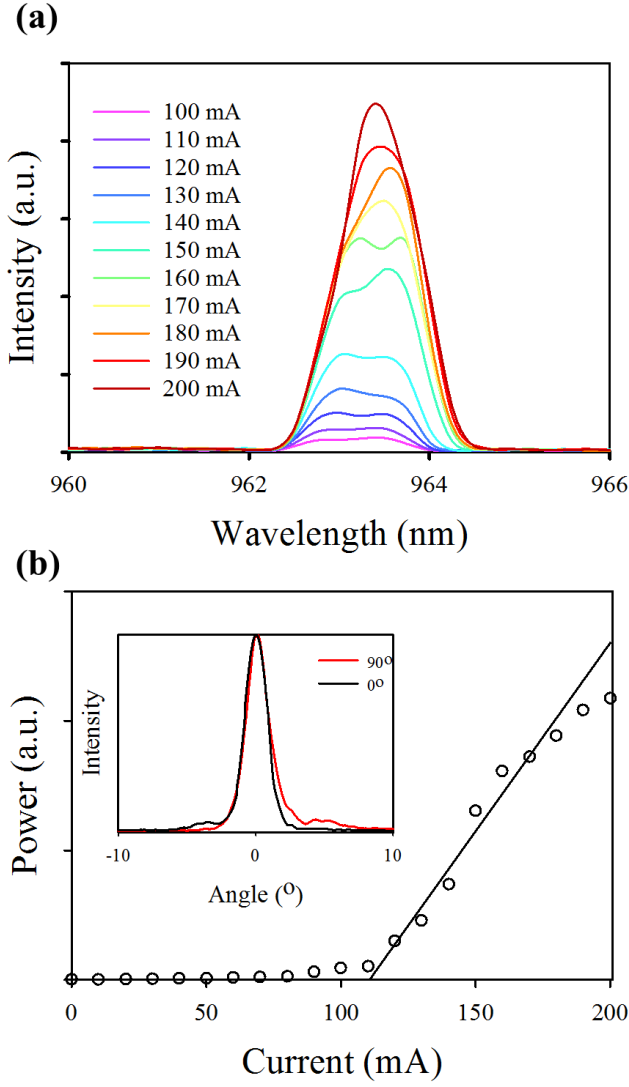


Figure 9 a) continuous wave EL spectra, b) LI of PCSEL device

Figure 10 shows the CW room-temperature LI of the virgin (black) PCSEL device, and with 1 (red), and 2 (blue) cleaves. With no cleave, the device has a threshold current of 112 mA ($J=1.43 \text{ kA}\cdot\text{cm}^{-2}$), the threshold is reduced to 88 mA ($J=1.12 \text{ kA}\cdot\text{cm}^{-2}$) when a single cleave is introduced to the PCSEL. The introduction of a second cleave does not reduce the threshold further, indicating that the lowest threshold 2D mode is selected by introducing external feedback to one side of the device. With additional feedback (one and 2 cleaves) an increased slope efficiency is also observed. Images show a 3D schematic representation of a PCSEL with no cleave, 1 cleave, and 2 cleaves.

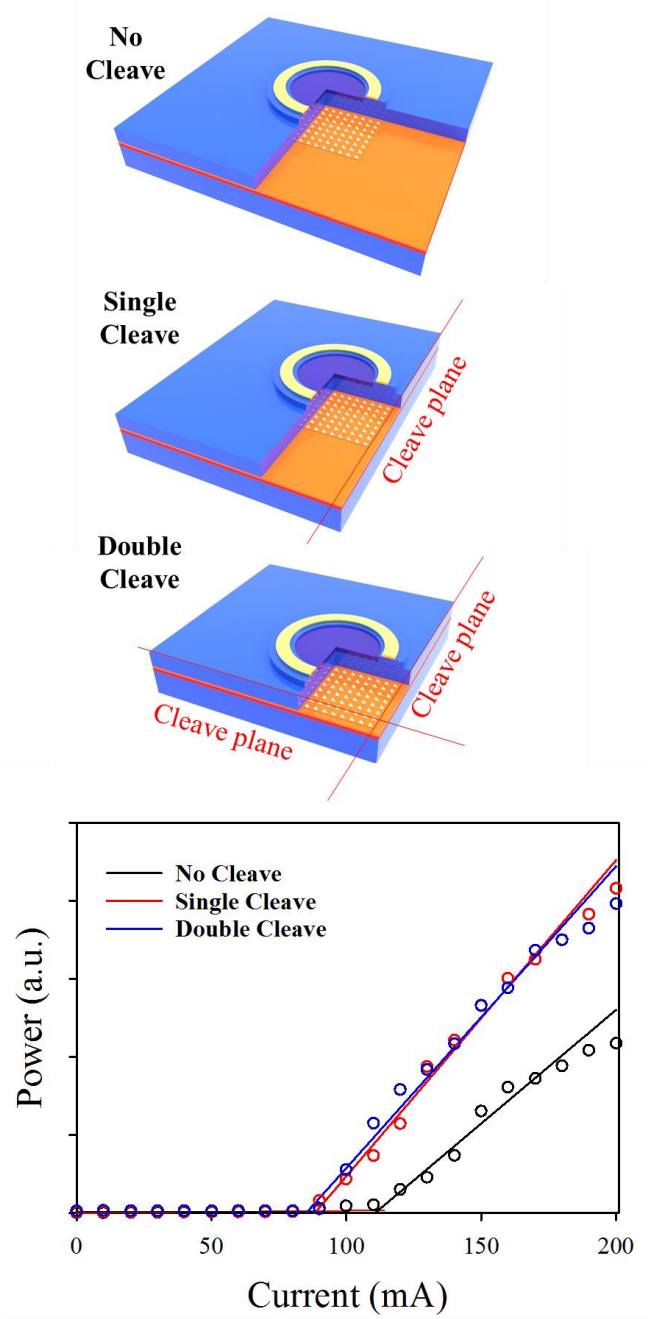


Figure 10 LI of the virgin (black) PCSEL device, and with 1 (red) and 2 (blue) cleaves.

Figure 11 shows the far-field at 1.11 μm for a PCSEL with no cleave (a), one cleave (b), and two cleaves (c). As described previously, for the virgin device, the far field pattern is shown to be symmetric, with a divergence of $\sim 1.2^\circ$. With the addition of a single cleaved facet, the far field pattern of the device becomes asymmetric with divergence $\sim 4^\circ$ for the 90° (perpendicular to cleave) direction and $\sim 1^\circ$ for the 0° (parallel to cleave) direction. The addition of feedback clearly modifies the field distribution within the device. The addition of a second cleave results in a divergence of $0.8\text{--}1^\circ$ indicating lasing is now taking place over a larger area of the PCSEL as compared to the virgin device.

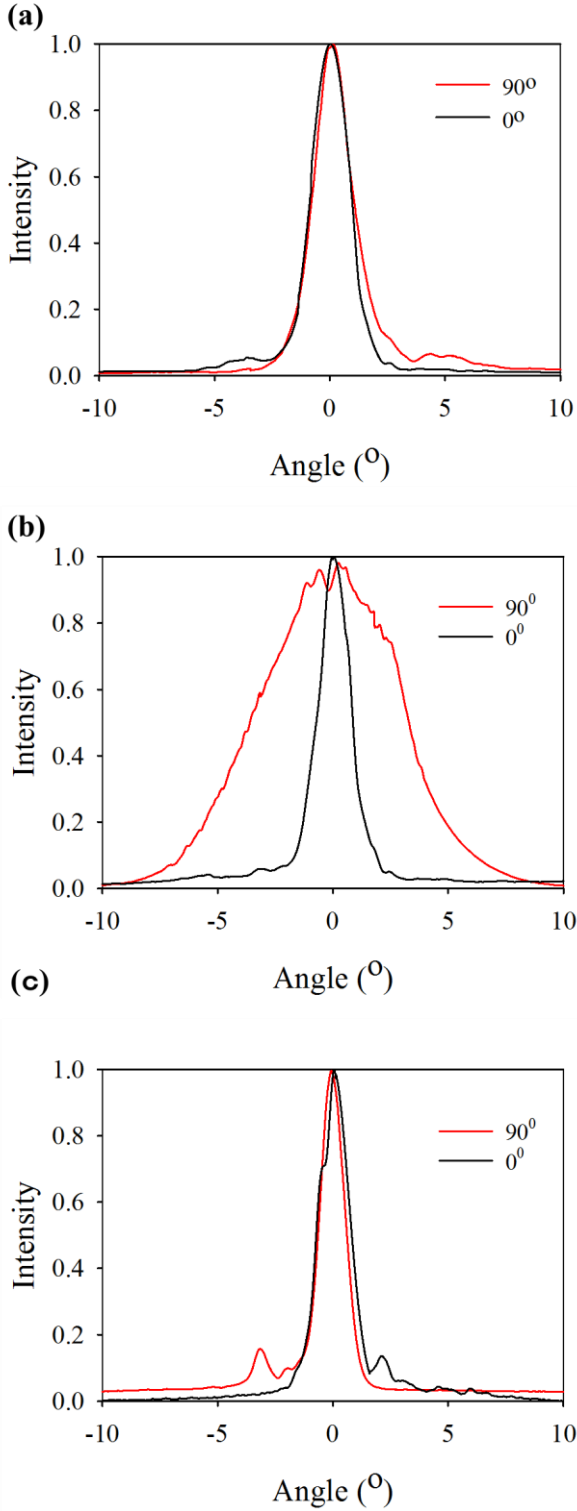


Figure 11 Far field pattern of PCSEL device (a) as fabricated, (b) with one cleave and, (c) with 2 cleaves

Figure 12 shows the CW room temperature EL spectra of the PCSEL device with no cleave (blue), 1 cleave (red), and 2 cleaves (black). As shown previously, the virgin device shows dual peak nature with peaks at 963 nm and 963.5 nm, with one cleave the spectra contains a single lasing peak at 963.75 nm. When a second cleave is introduced the peak wavelength

increases to 964.2 nm we attribute this to self-heating of the device caused by the 2 cleaves limiting the available area for heat dissipation to occur. However, we note that a shift in wavelength is an expected effect of lateral feedback [16].

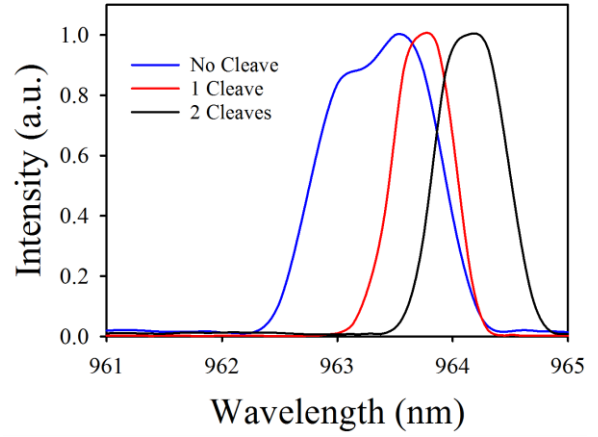


Figure 12 CW room temperature EL spectra for PCSEL device with no cleave (blue), 1 cleave (red), and 2 cleaves (black).

Figure 13 shows the peak wavelength as a function of current for a PCSEL device with no cleave, 1 cleave, and 2 cleaves. In each case the peak wavelength slightly increases as the applied current increases. For the virgin device two lasing modes of ~963 nm and ~963.5 nm are observed. The two distinct modes are attributed to modes C and D in Figure 2. The introduction of a single cleave results in single-mode behavior with a red-shift of the lasing peak of ~0.25 nm. Single mode operation with a red-shift of ~0.65 nm is observed for the case of two cleaves. For both one and two cleaves, single mode operation is maintained over all measured drive currents (up to 200 mA ($2.55 \text{ kA}\cdot\text{cm}^{-2}$)). As we have drastically changed the heatsinking of the devices we cannot rule out self-heating to be the cause of this redshift (~7 and 17° increase in junction temperature for the single and double cleaved device, respectively [10]).

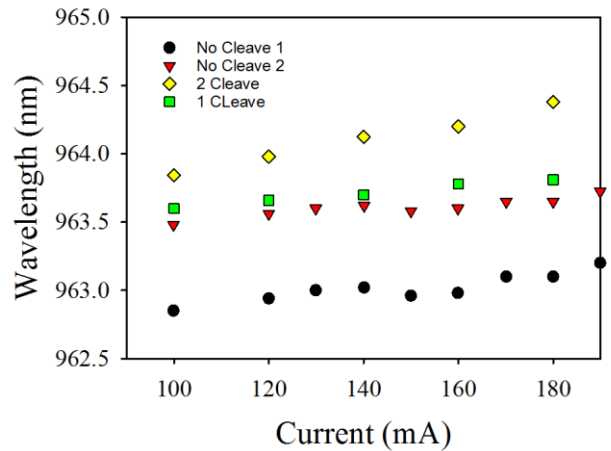


Figure 13 peak wavelength as a function of current for a PCSEL structure with no cleave (black/red), 1 cleave (green), 2 cleaves (red).

It is clear that mode selection is possible through the use of lateral feedback. As we have a distributed feedback phase, competition would be expected not necessarily to the next mode, but to slightly different feedback phases of the same mode. In the case of strong feedback ($R=1$), there is a significant change in lasing wavelength with varying feedback phase [16]. Spectral broadening of the lasing wavelength with increasing drive current may therefore be expected. However, recent work on spatial coherence effects in PCSELS suggest that injection locking occurs for large area PCSELS with a spatial variation of PC shape/period [22] and for individual PCSELS making up coupled arrays [5,22,23]. Further theoretical and experimental work is required to explore the mode stability under high power operation.

In order to harness feedback effects in a scalable manner, the PCSEL would need to be terminated with a DBR boundary. The incorporation of an electrically contacted coupling waveguide region [5, 24, 25] allows us to envisage electronic control of the feedback amplitude and phase, allowing all-electronic control of lasing position and beam divergence [24, 25].

V. CONCLUSIONS

In summary, an all semiconductor PCSEL based on epitaxial overgrowth with a square lattice of triangles has been studied theoretically and experimentally under a range of feedback regimes. Feedback to a single side of the device results in the selection of the lowest threshold 2D lasing mode and a modification to the field distribution within the device resulting in a broadening to the far-field in one direction. For an originally dual-mode laser, single mode operation is obtained over a larger area through the introduction of feedback to two facets.

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